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Uniqueness and Lipschitz stability of an inverse boundary value problem for time-harmonic elastic waves

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Abstract

We consider the inverse problem of determining the Lamé parameters and the density of a three-dimensional elastic body from the local time-harmonic Dirichlet-to-Neumann map. We prove uniqueness and Lipschitz stability of this inverse problem when the Lamé parameters and the density are assumed to be piecewise constant on a given domain partition.

Keywords: inverse boundary value problem, uniqueness, Lipschitz stability, time-harmonic elastic waves

(Some figures may appear in colour only in the online journal)

1. Introduction

We study the inverse boundary value problem for time-harmonic elastic waves. We consider isotropic elasticity, and allow partial boundary data. The Lamé parameters and the density are assumed to be piecewise constants on a given partitioning of the domain. The system of equations describing time-harmonic elastic waves is given by,

$$\begin{cases} \operatorname{div}(\mathbb{C}\hat{\nabla}u) + \rho\omega^2 u = 0 & \text{in } \Omega \subset \mathbb{R}^3, \\ u = \psi & \text{on } \partial\Omega, \end{cases} \tag{1}$$

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where Ω is an open and bounded domain with smooth boundary, $\hat{\nabla} u$ denotes the strain tensor, $\hat{\nabla} u := \frac{1}{2}(\nabla u + (\nabla u)^T)$, $\psi \in H^{1/2}(\partial\Omega)$ is the boundary displacement or source, and $\mathbb{C} \in L^\infty(\Omega)$ denotes the isotropic elasticity tensor with Lamé parameters λ, μ:

$$\mathbb{C} = \lambda I_3 \otimes I_3 + 2\mu \mathbb{I}_{\mathrm{sym}}, \text{ a.e. in } \Omega,$$

where I_3 is 3×3 identity matrix and $\mathbb{I}_{\mathrm{sym}}$ is the fourth order tensor such that $\mathbb{I}_{\mathrm{sym}} A = \hat{A}$, $\rho \in L^\infty(\Omega)$ is the density, and ω is the frequency. Here, we make use of the following notation for matrices and tensors: for 3×3 matrices A and B we set $A : B = \sum_{i,j=1}^{3} A_{ij} B_{ij}$ and $\hat{A} = \frac{1}{2}(A + A^T)$. We assume that

$$0 < \alpha_0 \leqslant \mu \leqslant \alpha_0^{-1}, \; \lambda \leqslant \alpha_0^{-1}, \; 2\mu + 3\lambda \geqslant \beta_0 > 0 \text{ a.e. in } \Omega, \tag{2}$$

$$\gamma_0 \leqslant \rho \leqslant \gamma_0^{-1}. \tag{3}$$

The Dirichlet-to-Neumann map, $\Lambda_{\mathbb{C},\rho}$, is defined by

$$\Lambda_{\mathbb{C},\rho} : H^{1/2}(\partial\Omega) \ni \psi \to (\mathbb{C}\hat{\nabla} u)\nu|_{\partial\Omega} \in H^{-1/2}(\partial\Omega),$$

where ν is the outward unit normal to $\partial\Omega$. We consider the inverse problem:

$$\text{determine } \mathbb{C}, \rho \text{ from } \Lambda_{\mathbb{C},\rho}.$$

For the static case (that is, $\omega = 0$) of our problem, Imanuvilov and Yamamoto [IY] proved, in dimension two, a uniqueness result for C^{10} Lamé parameters. In dimension three, Nakamura and Uhlmann [NU] proved uniqueness assuming that the Lamé parameters are C^∞ and that μ is close to a positive constant. Eskin and Ralston [ER] proved a similar result. Global uniqueness of the inverse problem in dimension three assuming general Lamé parameters remains an open problem. One key difficulty here is that there are two metrics involved in the elastic tensor. Beretta *et al* proved the uniqueness when the Lamé parameters are assumed to be piecewise constant. They proved a Lipschitz stability when interfaces of subdomains contain flat parts [BFV]; later, they extended this result to non-flat interfaces [BFMRV]. Alessandrini *et al* [AdCMR] proved a logarithmic stabilty estimate for the inverse problem of identifying an inclusion, where constant Lamé parameters are different from the background ones.

The time-harmonic problem under our consideration has a more practical setting. The key application we have in mind is (reflection) seismology, where Lamé parameters and density need to be recovered from the Dirichlet-to-Neumann map. In actual seismic acquisition, raw vibroseis data are modeled by the local Neumann-to-Dirichlet map: the boundary values are given by the normal traction underneath the base plate of a vibroseis and are zero ('free surface') elsewhere, while the particle displacement (in fact, velocity) is measured by geophones located in a subset of the boundary (Earth's surface). The applied signal is essentially time-harmonic (suppressing the sweep); see [B, 2.52 and (2.53)]. (The displacement needs to be measured also underneath the base plate.) The results presented here do not only hold for the Dirichlet-to-Neumann map, but also for the Neumann-to-Dirichlet map as the data (requiring only minor modifications of the proofs).

We consider piecewise constant Lamé parameters and density of the form

$$\mathbb{C}(x) = \sum_{j=1}^{N} (\lambda_j I_3 \otimes I_3 + 2\mu_j \mathbb{I}_{\mathrm{sym}}) \chi_{D_j}(x), \quad \rho(x) = \sum_{j=1}^{N} \rho_j \chi_{D_j}(x),$$

where the D_j's, $j = 1, \cdots, N$ are known disjoint Lipschitz domains and $\lambda_j, \mu_j, \rho_j, j = 1, \cdots, N$ are unknown constants. We establish uniqueness of the above mentioned inverse boundary value problem. We actually derive a Lipschitz stability estimate, and the uniqueness follows immediately. The method of proof follows the ideas introduced by Alessandrini and Vessella [AV] in the study of electrical impedance tomography (EIT) problems. The counterpart for scalar waves, that is, the inverse boundary value problem for the Helmholtz equation, was analyzed by Beretta *et al* [BdHQ].

The existence and the 'blow up' behavior of singular solutions close to a flat discontinuity are utilized in our proof. The quantitative estimate of unique continuation for elliptic systems, which is derived from a three spheres inequality, play an essential role in the procedure. We directly prove a log-type stability estimate for the Lamé parameters and the density combined by alternatingly estimating them along a walkway of subdomains. Uniqueness then follows from the stability estimate. From the restriction that the parameters to be recovered lie in a finite-dimensional space, a Lipschitz stability estimate is obtained.

A key complication addressed in this paper is the multiparameter aspect of this inverse problem. For the acoustic waves modeled by the equation

$$\nabla \cdot (\gamma \nabla u) + q\omega^2 u = 0, \tag{4}$$

Nachman [N] proved the unique recovery of $\gamma \in C^2$ and $q \in L^\infty$ with Dirichlet-to-Neumann maps at two different admissible frequencies ω_1, ω_2. For the optical tomography problem, that is, recovering simultaneously $a > 0$ and $c > 0$ in the partial differential equation

$$-\nabla \cdot (a\nabla u) + cu = 0,$$

from all possible boundary Dirichlet and Neumann pairs, Arridge and Lionheart [AL] demonstrated the non-uniqueness for general a and c. However, when a is piecewise constant and c is piecewise analytic, Harrach [H] proved the uniqueness of this inverse problem. In this paper, we prove, for our problem, that recovering a higher order coefficient and a lower order coefficient jointly, that are assumed to be piecewise constant, only needs single frequency data also. If we assume γ, q to be piecewise constant in (4), we can establish the uniqueness with single frequency data, following the methods of proof in this paper.

With the conditional Lipschitz stability which we obtain here, we can invoke iterative methods with guaranteed convergence for local reconstruction, such as the nonlinear Landweber iteration [dHQS1] and the nonlinear projected steepest descent algorithm [dHQS2] (including a stopping criterion which allows inaccurate data). For a numerical realization, we refer to [BdHFS]. In reflection seismology, iterative methods for solving inverse problems, casting these into optimization problems, have been collectively referred to as Full Waveform Inversion (FWI) through the use of the adjoint state method. These methods were introduced in this field of application by Chavent [C], Lailly [L] and Tarantola & Valette [T, TV] albeit for scalar waves. An early study of stability in dimension one can be found in Bamberger *et al* [BCL]. Mora [M] developed the adjoint state formulation for the case of elastic waves and carried out computational experiments; Crase *et al* [CPNMT] then carried out applications to field data. Advantages of using time-harmonic data, following specific workflows, were initially pointed out by Pratt and collaborators [P–PW]; Bunks *et al* [BKB] developed an important insight in the use of strictly finite-frequency data. In recent years, there has been a significant effort in further developing and applying these approaches (with emphasis on iterative Gauss–Newton methods)—in the absence of a notion of (conditional) uniqueness, stability or convergence—often in combination with intuitive strategies for selecting parts of the data. In exploration seismology, we mention the work of Gélis *et al* [GVG], Choi [CMS], Brossier *et al* [BOV1, BOV2] and Xu & McMechan [XM]; in global seismology, we mention the work of Tromp *et al* [TTL] and Fichtner & Trampert [FT].

The paper is organized as follows: in section 2, we summarize the main results. In section 3, we construct the singular solutions and establish the unique continuation for the system describing time-harmonic elastic waves. We also prove the Fréchet differentiability of the forward map, $(\mathbb{C}, \rho) \to \Lambda_{\mathbb{C},\rho}$. In section 4, we prove the main result.

2. Main result

2.1. Direct problem

We summarize some results concerning the well-posedness of problem (1). For the proof, we follow the lines of [BdHQS].

Proposition 2.1. *Let Ω be a bounded Lipschitz domain in $\mathbb{R}^3$, $f \in H^{-1}(\Omega)$ and $g \in H^{1/2}(\partial\Omega)$. Assume that λ, μ, ρ satisfy (2) and (3). Let λ_1^0 be the smallest Dirichlet eigenvalue of the operator $-\mathrm{div}(\mathbb{C}_0 \hat{\nabla} u)$ in Ω, where $\mathbb{C}_0 = \frac{\beta_0 - 3\alpha_0}{2} I_3 \otimes I_3 + 2\alpha_0 \mathbb{I}_{\mathrm{sym}}$. Then, for any $\omega^2 \in (0, \frac{\gamma_0 \lambda_1^0}{2}]$, there exists a unique solution of*

$$\begin{cases} \mathrm{div}(\mathbb{C}\hat{\nabla} u) + \rho\omega^2 u = f & \text{in } \Omega \subset \mathbb{R}^3, \\ u = g & \text{on } \partial\Omega, \end{cases} \tag{5}$$

satisfying

$$\|u\|_{H^1(\Omega)} \leqslant C(\|g\|_{H^{1/2}(\partial\Omega)} + \|f\|_{H^{-1}(\Omega)}), \tag{6}$$

where C depends on α_0, β_0, γ_0 and λ_1^0.

Proof. Without loss of generality, we let $g = 0$. Indeed, we can always introduce a $w = u - \tilde{g}$ where $\tilde{g} \in H^1(\Omega)$ is such that $\tilde{g} = g$ on $\partial\Omega$, which satisfies (5) with $g = 0$. We recall that

$$\lambda_1^0 = \min\left\{ \int_\Omega \mathbb{C}_0 \hat{\nabla} u : \hat{\nabla} u \,\middle|\, u \in H^1(\Omega),\ \|u\|_{L^2(\Omega)} = 1 \right\}, \tag{7}$$

and observe that $\mathbb{C} \geqslant \mathbb{C}_0$, that is, $(\mathbb{C} - \mathbb{C}_0)\hat{A} : \hat{A} \geqslant 0$ for any 3×3 matrix A.

We consider on $H_0^1(\Omega)$ the bilinear form

$$a(u, v) = \int_\Omega \mathbb{C}\hat{\nabla} u : \hat{\nabla} v \mathrm{d}x - \int_\Omega \omega^2 \rho u \cdot v \mathrm{d}x.$$

Then we can write problem (5) (for $g = 0$) in the weak form,

$$a(u, v) = -\langle f, v \rangle \quad \forall v \in H_0^1(\Omega).$$

Clearly $a(\cdot, \cdot)$ is continuous. We check now that $a(\cdot, \cdot)$ is coercive. To this aim, we recall the Korn inequality

$$\int_\Omega |\hat{\nabla} u|^2 \mathrm{d}x \leqslant 2 \int_\Omega |\nabla u|^2 \mathrm{d}x \tag{8}$$

for any $u \in H_0^1(\Omega)$ (using the matrix norm, $|A|^2 = A : A$ for any 3×3 matrix A). Furthermore,

$$\begin{aligned} a(u,u) &= \int_\Omega \mathbb{C}\hat{\nabla}u : \hat{\nabla}u\mathrm{d}x - \int_\Omega \omega^2\rho|u|^2\,\mathrm{d}x \\ &\geqslant \int_\Omega \mathbb{C}_0\hat{\nabla}u : \hat{\nabla}u\mathrm{d}x - \omega^2\gamma_0^{-1}\int_\Omega |u|^2\,\mathrm{d}x \\ &= \frac{1}{2}\int_\Omega \mathbb{C}_0\hat{\nabla}u : \hat{\nabla}u\mathrm{d}x + \frac{1}{2}\left\{\int_\Omega \mathbb{C}_0\hat{\nabla}u : \hat{\nabla}u\mathrm{d}x - 2\omega^2\gamma_0^{-1}\int_\Omega |u|^2\,\mathrm{d}x\right\}. \end{aligned}$$

By (7), the strong convexity of $\mathbb{C}_0$, the Korn inequality (8) and the Poincaré inequality, we have

$$\begin{aligned} a(u,u) &\geqslant \frac{\xi_0}{4}\int_\Omega |\nabla u|^2\,\mathrm{d}x + \frac{1}{2}\left\{\int_\Omega \mathbb{C}_0\hat{\nabla}u : \hat{\nabla}u\mathrm{d}x - 2\omega^2\gamma_0^{-1}\int_\Omega |u|^2\,\mathrm{d}x\right\} \\ &\geqslant \frac{\xi_0 C_P}{4}\|u\|^2_{H^1(\Omega)} \end{aligned}$$

indeed, where ξ_0 depends on α_0 and β_0 only and C_P is the Poincaré constant of Ω. By the Lax–Milgram lemma there exists a unique solution $u \in H^1_0(\Omega)$ to problem (5), and (6) holds. $\square$

Remark 2.2. We note that whenever ω is not in a particular countable subset of real numbers (the set of eigenfrequencies), problem (5) has a unique solution and estimate (6) holds with the constant C depending also on ω.

We let Σ be an open portion of $\partial\Omega$. We denote by $H^{1/2}_{co}(\Sigma)$ the space

$$H^{1/2}_{co}(\Sigma) := \{\phi \in H^{1/2}(\partial\Omega) \mid \mathrm{supp}\,\phi \Subset \Sigma\}$$

and by $H^{-1/2}_{co}(\Sigma)$ the topological dual of $H^{1/2}_{co}(\Sigma)$. We denote by $\langle\cdot,\cdot\rangle$ the dual pairing between $H^{1/2}_{co}(\Sigma)$ and $H^{-1/2}_{co}(\Sigma)$ based on the $L^2(\Sigma)$ inner product. By proposition 2.1 it follows that for any $\psi \in H^{1/2}_{co}(\Sigma)$ there exists a unique vector-valued function $u \in H^1(\Omega)$ that is a weak solution of the Dirichlet problem (1). We define the local Dirichlet-to-Neumann map $\Lambda^\Sigma_{\mathbb{C},\rho}$ as

$$\Lambda^\Sigma_{\mathbb{C},\rho} : H^{1/2}_{co}(\Sigma) \ni \psi \to (\mathbb{C}\hat{\nabla}u)\nu|_\Sigma \in H^{-1/2}_{co}(\Sigma).$$

We have $\Lambda_{\mathbb{C},\rho} = \Lambda^{\partial\Omega}_{\mathbb{C},\rho}$. The map $\Lambda^\Sigma_{\mathbb{C},\rho}$ can be identified with the bilinear form on $H^{1/2}_{co}(\Sigma) \times H^{-1/2}_{co}(\Sigma)$,

$$\hat{\Lambda}^\Sigma_{\mathbb{C},\rho}(\psi,\phi) := \langle\Lambda^\Sigma_{\mathbb{C},\rho}\psi,\phi\rangle = \int_\Omega (\mathbb{C}\hat{\nabla}u : \hat{\nabla}v - \rho\omega^2 u\cdot v)\mathrm{d}x, \tag{9}$$

for all $\psi,\phi \in H^{1/2}_{co}(\Sigma)$, where u solves (1) and v is any $H^1(\Omega)$ function such that $v = \phi$ on $\partial\Omega$. We shall denote by $\|\cdot\|_\star$ the norm in $\mathcal{L}(H^{1/2}_{co}(\Sigma), H^{-1/2}_{co}(\Sigma))$ defined by

$$\|T\|_\star = \sup\{\langle T\psi,\phi\rangle \mid \psi,\phi \in H^{1/2}_{co}(\Sigma), \|\psi\|_{H^{1/2}_{co}(\Sigma)} = \|\phi\|_{H^{1/2}_{co}(\Sigma)} = 1\}.$$

2.2. Notation and definitions

For every $x \in \mathbb{R}^3$ we set $x = (x', x_3)$ where $x' \in \mathbb{R}^2$ and $x_3 \in \mathbb{R}$. For every $x \in \mathbb{R}^3$, r and L positive real numbers we denote by $B_r(x)$, $B'_r(x')$ and $Q_{r,L}$ the open ball in $\mathbb{R}^3$ centered at x of radius r, the open ball in $\mathbb{R}^2$ centered at x' of radius r and the cylinder $B'_r(x') \times (x_3 - Lr, x_3 + Lr)$, respectively; $B_r(0)$, $B'_r(0)$ and $Q_{r,L}(0)$ will be denoted by B_r, B'_r and $Q_{r,L}$, respectively. We will also write $\mathbb{R}^3_+ = \{(x', x_3) \in \mathbb{R}^3 : x_3 > 0\}$, $\mathbb{R}^3_- = \{(x', x_3) \in \mathbb{R}^3 : x_3 < 0\}$, $B_r^+ = B_r \cap \mathbb{R}^3_+$, and $B_r^- = B_r \cap \mathbb{R}^3_-$. For any subset D of $\mathbb{R}^3$ and any $h > 0$, we let

$$(D)_h = \{x \in D \mid \operatorname{dist}(x, \mathbb{R}^3 \setminus D) > h\}.$$

Definition 2.3. Let Ω be a bounded domain in $\mathbb{R}^3$. We say that a portion $\Sigma \subset \partial\Omega$ is of Lipschitz class with constants $r_0 > 0, L \geqslant 1$ if for any point $P \in \Sigma$, there exists a rigid transformation of coordinates under which $P = 0$ and

$$\Omega \cap Q_{r_0,L} = \{(x', x_3) \in Q_{r_0,L} \mid x_3 > \psi(x')\},$$

where ψ is a Lipschitz continuous function in B'_{r_0} such that

$$\psi(0) = 0 \text{ and } \|\psi\|_{C^{0,1}(B'_{r_0})} \leqslant Lr_0.$$

We say that Ω is of Lipschitz class with constants r_0 and L if $\partial\Omega$ is of Lipschitz class with the same constants.

2.3. Main assumptions

Let $A, L, \alpha_0, \beta_0, \gamma_0, N$ be given positive numbers such that $N \in \mathbb{N}$, $\alpha_0 \in (0, 1)$, $\beta_0 \in (0, 2)$, $\gamma_0 \in (0, 1)$ and $L > 1$. We shall refer to them as the prior data.

In the sequel we will introduce a various constants that we will always denote by C. The values of these constants might differ from one another, but we will always have $C > 1$.

Assumption 2.4. ([BFV]). The domain $\Omega \subset \mathbb{R}^3$ is open and bounded with

$$|\Omega| \leqslant A,$$

and

$$\bar{\Omega} = \cup_{j=1}^N \bar{D}_j,$$

where $D_j, j = 1, \ldots, N$ are connected and pairwise non-overlapping open subdomains of Lipschitz class with constants 1,L. Moreover, there exists a region, say D_1, such that $\partial D_1 \cap \partial\Omega$ contains an open flat part, Σ, and that for every $j \in \{2, \ldots, N\}$ there exist $j_1, \ldots, j_M \in \{1, \ldots, N\}$ such that

$$D_{j_1} = D_1, \quad D_{j_M} = D_j$$

and, for every $k = 2, \ldots, M$

$$\partial D_{j_{k-1}} \cap \partial D_{j_k}$$

contains a flat portion Σ_k such that

$$\Sigma_k \Subset \Omega, \text{ for all } k = 2, \ldots, M.$$

Furthermore, for $k = 1, \ldots, M$, there exists $P_k \in \Sigma_k$ and a rigid transformation of coordinates such that $P_k = 0$ and

$$\Sigma_k \cap Q_{1/3,L} = \{x \in Q_{1/3,L} : x_3 = 0\},$$

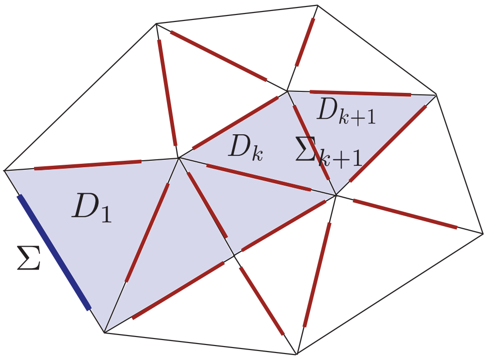

Figure 1. A domain partition including D_1.

$$D_{j_k} \cap Q_{1/3,L} = \{x \in Q_{1/3,L} : x_3 < 0\},$$

$$D_{j_{k-1}} \cap Q_{1/3,L} = \{x \in Q_{1/3,L} : x_3 > 0\};$$

here, we set $\Sigma_1 = \Sigma$. We will refer to $D_{j_1}, \ldots, D_{j_M}$ as a chain of subdomains connecting D_1 to D_j. For any $k \in \{1, \ldots, M\}$ we will denote by n_k the exterior unit vector to ∂D_k at P_k.

An example of such a domain partition with Lipschitz class subdomains is an unstructured tetrahedral mesh as shown in figure 1.

Assumption 2.5. The stiffness tensor, $\mathbb{C}$, is isotropic and piecewise constant, that is,

$$\mathbb{C} = \sum_{j=1}^{N} \mathbb{C}_j \chi_{D_j}(x), \quad \mathbb{C}_j = \lambda_j I_3 \otimes I_3 + 2\mu_j \mathbb{I}_{\mathrm{sym}},$$

where the constants λ_j and μ_j satisfy (see (2))

$$0 < \alpha_0 \leqslant \mu_j \leqslant \alpha_0^{-1}, \quad \lambda_j \leqslant \alpha_0^{-1}, \quad 2\mu_j + 3\lambda_j \geqslant \beta_0 > 0,\ j = 1, \ldots, N. \tag{10}$$

The density, ρ, is of the form,

$$\rho = \sum_{j=1}^{N} \rho_j \chi_{D_j}(x),$$

where the constants ρ_j satisfy (see (3))

$$\gamma_0 \leqslant \rho_j \leqslant \gamma_0^{-1},\ j = 1, \ldots, N.$$

Assumption 2.6. Let λ_1^0 be the smallest Dirichlet eigenvalue of operator $-\mathrm{div}(\mathbb{C}_0 \hat{\nabla} u)$ in Ω as before,

$$\omega^2 \leqslant \frac{\gamma_0 \lambda_1^0}{2}.$$

2.4. Statement of the main result

We define for any set $D \in \mathbb{R}^3$,

$$d_D((\mathbb{C}^1, \rho^1), (\mathbb{C}^2, \rho^2)) = \max\{\|\lambda^1 - \lambda^2\|_{L^\infty(D)}, \|\mu^1 - \mu^2\|_{L^\infty(D)}, \|\rho^1 - \rho^2\|_{L^\infty(D)}\}.$$

Theorem 2.7. *Let $(\mathbb{C}^{1,2}, \rho^{1,2})$ satisfy assumption 2.5. Let Ω and Σ satisfy assumption 2.4 and ω satisfy assumption 2.6. If $\Lambda^\Sigma_{\mathbb{C}^2,\rho^2} = \Lambda^\Sigma_{\mathbb{C}^1,\rho^1}$ then $\mathbb{C}^1 = \mathbb{C}^2$ and $\rho^1 = \rho^2$. Moreover, there exists a positive constant C depending on L, A, N, α_0, β_0, γ_0 and λ_1^0 only, such that*

$$d_\Omega((\mathbb{C}^1, \rho^1), (\mathbb{C}^2, \rho^2)) \leqslant C\, \|\Lambda^\Sigma_{\mathbb{C}^1,\rho^1} - \Lambda^\Sigma_{\mathbb{C}^2,\rho^2}\|_\star. \tag{11}$$

In preparation of the proof, we introduce the forward map associated with the inverse problem. We let $\underline{L} := (\lambda_1, \ldots, \lambda_N, \mu_1, \ldots, \mu_N, \rho_1, \ldots, \rho_N)$ denote a vector in $\mathbb{R}^{3N}$ and $\mathcal{A}$ stand for the open subset of $\mathbb{R}^{3N}$ defined by

$$\mathcal{A} := \left\{\underline{L} \in \mathbb{R}^{3N} \;\middle|\; \frac{\alpha_0}{2} < \mu_j < \frac{2}{\alpha_0}, \lambda_j < \frac{2}{\alpha_0}, 2\mu_j + 3\lambda_j > \frac{\beta_0}{2}, \frac{\gamma_0}{2} < \rho_j < \frac{2}{\gamma_0}, j = 1, \ldots, N\right\}. \tag{12}$$

For each vector $\underline{L} \in \mathcal{A}$ we can define a piecewise constant stiffness tensor $\mathbb{C}_{\underline{L}}$, and a density $\rho_{\underline{L}}$, with

$$\|\underline{L}\|_\infty = \max_{j=1,\ldots,N}\{\sup\{|\lambda_j|, \mu_j, |\rho_j|\}\}.$$

The forward map is defined as

$$F:\ \mathcal{A} \to \mathcal{L}(H^{1/2}_{co}(\Sigma), H^{-1/2}_{co}(\Sigma)), \qquad \underline{L} \to F(\underline{L}) = \Lambda^\Sigma_{\mathbb{C}_{\underline{L}},\rho_{\underline{L}}}. \tag{13}$$

We can identify F with a map $\tilde{F} : \mathcal{A} \to \mathcal{B}$ upon identifying $\tilde{F}(\underline{L})$ with the bilinear form, $\tilde{\Lambda}^\Sigma_{\mathbb{C}_{\underline{L}},\rho_{\underline{L}}}$, on $H^{1/2}_{co}(\Sigma) \times H^{-1/2}_{co}(\Sigma)$ (see (9)); $\mathcal{B}$ is the Banach space of this bilinear form with the standard norm. In the sequel, we will write F and $\Lambda^\Sigma_{\mathbb{C}_{\underline{L}},\rho_{\underline{L}}}$ instead of $\tilde{F}$ and $\tilde{\Lambda}^\Sigma_{\mathbb{C}_{\underline{L}},\rho_{\underline{L}}}$. We denote

$$\mathbf{K} := \{\underline{L} \in \mathcal{A} \mid \alpha_0 \leqslant \mu_j \leqslant \alpha_0^{-1}, \lambda_j \leqslant \alpha_0^{-1}, 2\mu_j + 3\lambda_j \geqslant \beta_0, \gamma_0 \leqslant \rho_j \leqslant \gamma_0^{-1}, j = 1, \ldots, N\}.$$

Then the stability estimate in theorem 2.7 can be stated as follows:

$$\|\underline{L}^1 - \underline{L}^2\|_\infty \leqslant C\|F(\underline{L}^1) - F(\underline{L}^2)\|_\star,$$

for every $\underline{L}^1, \underline{L}^2$ in $\mathbf{K}$. We note that theorem 2.7 implies that F is injective and that its inverse is Lipschitz continuous.

Remark 2.8. Assumption 2.6 in theorem 2.7 can be relaxed to include any ω that is not in the set of eigenfrequencies. Then the constant C will also depend on the distance between ω and the set of eigenfrequencies.

Remark 2.9. We emphasize here that the Lipschitz constant C in the stability estimate (11) grows exponentially with N, the number of subdomains. For such behaviors of this type of inverse problems, we refer to [BdHQ] and [Ron].

3. Preliminary results

Here, we follow Beretta *et al* [BFMRV, BFV]. We summarize the relevant results in their work and adapt them to the time-harmonic problem. We begin this section with Alessandrini's identity [A, I]. We let u_k be solutions to

$$\operatorname{div}(\mathbb{C}^k\hat{\nabla}u_k) + \rho^k\omega^2 u_k = 0 \quad \text{in } \Omega$$

for $k = 1, 2$, where $\mathbb{C}^k$, ρ^k satisfy assumption 2.5. Then

$$\int_\Omega ((\mathbb{C}^1 - \mathbb{C}^2)\hat{\nabla}u_1 : \hat{\nabla}u_2 - (\rho^1 - \rho^2)\omega^2 u_1 \cdot u_2)\mathrm{d}x = \langle(\Lambda_{\mathbb{C}^1,\rho^1} - \Lambda_{\mathbb{C}^2,\rho^2})u_1, u_2\rangle. \tag{14}$$

3.1. Fréchet differentiability of F

Here, we prove the Fréchet differentiability of the forward map, F.

Proposition 3.1. *Under assumptions 2.4–2.6, the map*

$$F : \mathcal{A} \to \mathcal{L}(H^{1/2}_{co}(\Sigma), H^{-1/2}_{co}(\Sigma))$$

is Frechét differentiable in $\mathcal{A}$ *and*

$$\langle DF(\underline{L})[\underline{H}]\psi, \phi\rangle = \int_\Omega (\mathbb{H}\hat{\nabla}u_{\underline{L}} : \hat{\nabla}v_{\underline{L}} - h\omega^2 u_{\underline{L}} \cdot v_{\underline{L}})\mathrm{d}x, \tag{15}$$

where $\mathbb{H} = \mathbb{C}_{\underline{H}}, h = \rho_{\underline{H}}$. *Moreover,* $DF : \mathcal{A} \to \mathcal{L}(\mathbb{R}^{3N}, \mathcal{L}(H^{1/2}_{co}(\Sigma), H^{-1/2}_{co}(\Sigma)))$ *is Lipschitz continuous with Lipschitz constant* C_{DF} *depending on* A, L, α_0, β_0, γ_0, λ_1^0 *only.*

Proof. Fix $\underline{L} \in \mathcal{A}$ and let $\underline{H} \in \mathbb{R}^{3N}$ such that $\|\underline{H}\|_\infty$ is sufficiently small. By (14) we have

$$\langle(F(\underline{L} + \underline{H}) - F(\underline{L}))\psi, \phi\rangle = \int_\Omega \mathbb{H}\hat{\nabla}u_{\underline{L}+\underline{H}} : \hat{\nabla}v_{\underline{L}}\mathrm{d}x - \int_\Omega h\omega^2 u_{\underline{L}+\underline{H}} \cdot v_{\underline{L}}\mathrm{d}x.$$

Hence, by setting

$$\begin{aligned}\eta &:= \langle(F(\underline{L} + \underline{H}) - F(\underline{L}))\psi, \phi\rangle - \int_\Omega \mathbb{H}\hat{\nabla}u_{\underline{L}} : \hat{\nabla}v_{\underline{L}}\mathrm{d}x + \int_\Omega h\omega^2 u_{\underline{L}} \cdot v_{\underline{L}}\mathrm{d}x \\ &= \int_\Omega \mathbb{H}\hat{\nabla}(u_{\underline{L}+\underline{H}} - u_{\underline{L}}) : \hat{\nabla}v_{\underline{L}}\mathrm{d}x - \int_\Omega h\omega^2(u_{\underline{L}+\underline{H}} - u_{\underline{L}}) \cdot v_{\underline{L}}\mathrm{d}x,\end{aligned} \tag{16}$$

we find that

$$|\eta| \leqslant C\|\underline{H}\|_\infty \|\nabla(u_{\underline{L}+\underline{H}} - u_{\underline{L}})\|_{L^2(\Omega)} \|\phi\|_{H^{1/2}_{co}(\Sigma)}, \tag{17}$$

where C depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0$ only. We estimate $\|\nabla(u_{\underline{L}+\underline{H}} - u_{\underline{L}})\|_{L^2(\Omega)}$. We observe that $w := u_{\underline{L}+\underline{H}} - u_{\underline{L}}$ is the solution to

$$\begin{cases}\operatorname{div}(\mathbb{C}_{\underline{L}}\hat{\nabla}w) + \rho\omega^2 w = -\operatorname{div}(\mathbb{H}\hat{\nabla}u_{\underline{L}+\underline{H}}) - h\omega^2 u_{\underline{L}+\underline{H}} & \text{in } \Omega, \\ w = 0 & \text{on } \partial\Omega.\end{cases} \tag{18}$$

By proposition 2.1, we have

$$\begin{aligned}\|\nabla w\|_{L^2(\Omega)} &\leqslant C\|w\|_{H^1(\Omega)}\\ &\leqslant C\|\mathrm{div}\,(\mathbb{H}\hat{\nabla}u_{\underline{L}+\underline{H}})\|_{H^{-1}(\Omega)}+C\|h\omega^2 u_{\underline{L}+\underline{H}}\|_{H^{-1}(\Omega)}\\ &\leqslant C\|\mathbb{H}\hat{\nabla}u_{\underline{L}+\underline{H}}\|_{L^2(\Omega)}+C\|h\omega^2 u_{\underline{L}+\underline{H}}\|_{H^{-1}(\Omega)}\\ &\leqslant C\|\underline{H}\|_\infty\|u_{\underline{L}+\underline{H}}\|_{H^1(\Omega)}+C\|\underline{H}\|_\infty\|u_{\underline{L}+\underline{H}}\|_{L^2(\Omega)}\\ &\leqslant C\|\underline{H}\|_\infty\|\psi\|_{H^{1/2}_{co}(\Sigma)},\end{aligned}\tag{19}$$

where C depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0$. By inserting (19) into (17) we get

$$|\eta|\leqslant C\|\underline{H}\|_\infty^2\|\psi\|_{H^{1/2}_{co}(\Sigma)}\|\phi\|_{H^{1/2}_{co}(\Sigma)},\tag{20}$$

that yields (15).

We now prove the Lipschitz continuity of DF. Let $\underline{L}^1, \underline{L}^2\in\mathcal{A}$ and set

$$\begin{aligned}\xi &:= \langle (DF(\underline{L}^2)-DF(\underline{L}^1))[\underline{H}]\psi,\phi\rangle\\ &= \int_\Omega(\mathbb{H}\hat{\nabla}u_{\underline{L}^2}:v_{\underline{L}^2}-\mathbb{H}\hat{\nabla}u_{\underline{L}^1}:v_{\underline{L}^1})\mathrm{d}x+\int_\Omega(h\omega^2u_{\underline{L}^2}\cdot v_{\underline{L}^2}-h\omega^2u_{\underline{L}^1}\cdot v_{\underline{L}^1})\mathrm{d}x\\ &= \int_\Omega\mathbb{H}(\hat{\nabla}u_{\underline{L}^2}-\hat{\nabla}u_{\underline{L}^1}):\hat{\nabla}v_{\underline{L}^2}\mathrm{d}x+\int_\Omega\mathbb{H}\hat{\nabla}u_{\underline{L}^1}:(\hat{\nabla}v_{\underline{L}^2}-\hat{\nabla}v_{\underline{L}^1})\mathrm{d}x\\ &\quad+\int_\Omega h\omega^2(u_{\underline{L}^2}-u_{\underline{L}^1})\cdot v_{\underline{L}^2}\mathrm{d}x+\int_\Omega h\omega^2u_{\underline{L}^1}\cdot(v_{\underline{L}^2}-v_{\underline{L}^1})\mathrm{d}x.\end{aligned}$$

By reasoning as we did to derive (20) we obtain

$$|\xi|\leqslant C_{DF}\|\underline{H}\|_\infty\|\underline{L}^2-\underline{L}^1\|_\infty\|\psi\|_{H^{1/2}_{co}(\Sigma)}\|\phi\|_{H^{1/2}_{co}(\Sigma)},$$

where C_{DF} depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0$. □

3.2. Further notation and definitions

Construction of an augmented domain and extension of $\mathbb{C}$ and ρ. First we extend the domain Ω to a new domain Ω_0 such that $\partial\Omega_0$ is of Lipschitz class and $B_{1/C}(P_1)\cap\Sigma\subset\Omega_0$, for some suitable constant $C\geqslant 1$ depending only on L. We proceed as in [ARRV]. We set

$$\eta_1 = 1/C_L,\ \text{where}\ C_L=\frac{3\sqrt{1+L^2}}{L},\tag{21}$$

and define, for every $x'\in B'_{\frac{1}{3}}$

$$\psi^+(x')=\begin{cases}\dfrac{\eta_1}{2} & \text{for } |x'|\leqslant\dfrac{\eta_1}{4L}\\ \eta_1-2L|x'| & \text{for } \dfrac{\eta_1}{4L}<|x'|\leqslant\dfrac{\eta_1}{2L}\\ 0 & \text{for } |x'|>\dfrac{\eta_1}{2L}.\end{cases}$$

We observe that for every $x' \in B'_{1/3}$, $|\psi^+(x')| \leqslant \frac{\eta_1}{2}$ and $|\nabla_{x'}\psi^+(x')| \leqslant 2L$. Next, we denote by

$$D_0 = \{x = (x', x_3) \in Q_{1/3,L} \mid 0 \leqslant x_3 < \psi^+(x')\},$$

$$\Omega_0 = \Omega \cup D_0.$$

We have

(i) Ω_0 has a Lipschitz boundary with constants $\frac{1}{3}, 3L$;

(ii)

$$\Omega_0 \supset Q_{1/4LC_L,L}.$$

Let $\mathbb{C}$ be an isotropic tensor that satisfies assumption 2.5. We extend $\mathbb{C}$ to Ω_0 such that $\mathbb{C}|_{D_0} = \mathbb{C}_0$. We also extend ρ such that $\rho|_{D_0} = 1$. Then $\mathbb{C}, \rho$ are of the form

$$\mathbb{C} = \sum_{j=0}^{N} \mathbb{C}_j \chi_{D_j}(x), \tag{22}$$

$$\rho = \sum_{j=0}^{N} \rho_j\, \chi_{D_j}(x). \tag{23}$$

Construction of a walkway. We fix $j \in \{1, \ldots, N\}$ and let $D_{j_1}, \ldots, D_{j_M}$ be a chain of domains connecting D_1 to D_j. We set $D_k = D_{j_k}$, $k = 1, \ldots, M$. By [ARRV] proposition 5.5, there exists $C'_L \geqslant 1$ depending on L only, such that $(D_k)_h$ is connected for every $k \in \{1, \ldots, M\}$ and every $h \in (0, 1/C'_L)$. We introduce

$$h_0 = \min\left\{\frac{1}{6}, \frac{1}{C'_L}, \frac{\eta_1}{8\sqrt{1+4L^2}}\right\} \tag{24}$$

where η_1 is as in (21).

Furthermore

(i) $Q_{(k)}$, $k = 1, \ldots, M$, is the cylinder centered at P_k such that by a rigid transformation of coordinates under which $P_k = 0$ and Σ_k belongs to the plane $\{(x', 0)\}$, and $Q_{(k)} = Q_{\eta_1/4L,L}$. We also denote $Q^-_{(M)} = Q_{(M)} \cap D_{M-1}$;

(ii) $\mathcal{K}$ is the interior part of the set $\bigcup_{k=1}^{M-1} \bar{D}_i$;

(iii) $\mathcal{K}_h = \cup_{k=1}^{M-1} (D_i)_h$, for every $h \in (0, h_0)$;

(iv)

$$\tilde{\mathcal{K}}_h = \mathcal{K}_h \cup Q^-_{(M)} \cup \bigcup_{k=1}^{M-1} Q_{(k)}; \tag{25}$$

(v)

$$K_0 = \left\{x \in D_0 \mid \mathrm{dist}(x, \partial\Omega) > \frac{\eta_1}{8}\right\}.$$

It is straightforward to verify that $\tilde{K}_h$ is connected and of Lipschitz class for every $h \in (0, h_0)$ and that

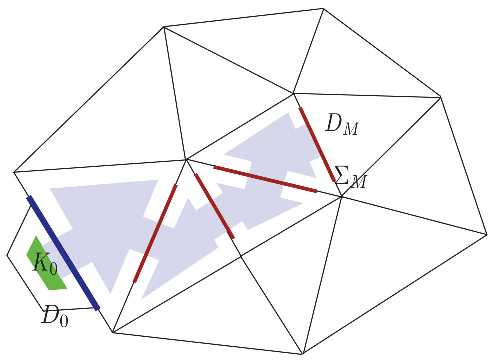

Figure 2. A path of the walkway.

$$K_0 \supset B'_{\eta_1/4L}(P_1) \times \left(\frac{\eta_1}{8}, \frac{\eta_1}{4}\right). \tag{26}$$

A path of the walkway is exhibited in Figure 2.

3.3. Existence of singular solutions

Next, we construct singular solutions to the system describing time-harmonic elastic waves. We prove the stability estimates for our inverse problems by studying the behavior of singular solutions.

3.3.1. Static fundamental solution in the biphase laminate. In order to construct singular solutions, we make use of special fundamental solutions constructed by Rongved [Rong] for isotropic biphase laminates. Consider

$$\mathbb{C}_b = \mathbb{C}^+ \chi_{\mathbb{R}^3_+} + \mathbb{C}^- \chi_{\mathbb{R}^3_-},$$

where $\mathbb{C}^+$ and $\mathbb{C}^-$ are constant isotropic stiffness tensors given by

$$\mathbb{C}^+ = \lambda I_3 \otimes I_3 + 2\mu \mathbb{I}_{\mathrm{sym}},\ \ \mathbb{C}^- = \lambda' I_3 \otimes I_3 + 2\mu' \mathbb{I}_{\mathrm{sym}},$$

with λ, μ and λ', μ' satisfying (10).

By [Rong], there exists a fundamental solution $\Gamma : \{(x, y) \mid x \in \mathbb{R}^3, y \in \mathbb{R}^3, x \neq y\} \to \mathbb{R}^{3\times 3}$ such that

$$\mathrm{div}(\mathbb{C}_b \hat{\nabla} \Gamma(\cdot, y)) = -\delta_y I_3.$$

Here δ_y is the Dirac distribution concentrated at y. We point out some properties of Γ. First of all, it is a fundamental solution, in the sense that $\Gamma(x, y)$ is continuous in $\{(x, y) \in \mathbb{R}^3 \times \mathbb{R}^3 \mid x \neq y\}$, $\Gamma(x, \cdot)$ is locally integrable in $\mathbb{R}^3$ for all $x \in \mathbb{R}^3$, and, for every vector valued function $\phi \in C_0^\infty(\mathbb{R}^3)$, we have

$$\int_{\mathbb{R}^3} \mathbb{C}_b \hat{\nabla} \Gamma(\cdot, y) : \hat{\nabla} \phi \, \mathrm{d}x = \phi(y).$$

Furthermore, for every $x, y \in \mathbb{R}^3, x \neq y$, we have

$$|\Gamma(x, y)| \leqslant \frac{C}{|x - y|}$$

and

$$|\nabla\Gamma(x,y)| \leqslant \frac{C}{|x-y|^2},$$

while for any $r > 0$,

$$\|\nabla\Gamma(\cdot,y)\|_{L^2(\mathbb{R}^3\setminus B_r(y))} \leqslant \frac{C}{r^{1/2}}, \tag{27}$$

where C depends on α_0, β_0 only.

3.3.2. Time-harmonic singular solutions. Let $\mathfrak{F}$ denote the union of the flats parts of $\bigcup_{j=1}^N \partial D_j$. Let $\mathcal{G} = \cup_{j=0}^N \partial D_j \setminus \mathfrak{F}$. Let $\mathbb{C} = \sum_{j=0}^N \mathbb{C}_j \chi_{D_j}$ where the tensors $\mathbb{C}_j$ satisfy assumption 2.5. Let $y \in \Omega_0 \setminus \mathcal{G}$ and let $r = \min(1/4, \text{ dist }(y, \mathcal{G} \cup \partial\Omega_0))$. Then, in the ball $B_r(y)$, either $\mathbb{C}$ is constant, $\mathbb{C} = \mathbb{C}_j$ or $\mathbb{C} = \mathbb{C}_j + (\mathbb{C}_{j+1} - \mathbb{C}_j)\chi_{\{x_3>a\}}$ for some a with $|a| < r$. We write

$$\mathbb{C}_y = \begin{cases} \mathbb{C}_j \ \text{ if } \mathbb{C} = \mathbb{C}_j \text{ in } B_r(y), \\ \mathbb{C}_j + (\mathbb{C}_{j+1} - \mathbb{C}_j)\chi_{\{x_3>a\}} \text{ otherwise,} \end{cases}$$

and consider the biphase fundamental solution satisfying

$$\text{div}(\mathbb{C}_y \hat{\nabla}\Gamma(\cdot,y)) = -\delta_y I_3 \text{ in } \mathbb{R}^3.$$

Proposition 3.2. *Let Ω_0, $\mathbb{C}$ and ω satisfy assumptions 2.4–2.6. Then, for $y \in \Omega_0 \setminus \mathcal{G}$, there exists only one function $G(\cdot,y)$, which is continuous in $\Omega \setminus \{y\}$, such that*

$$\int_{\Omega_0} (\mathbb{C}\hat{\nabla}G(\cdot,y) : \hat{\nabla}\phi - \rho\omega^2 G(\cdot,y)\cdot\phi)\mathrm{d}x = \phi(y),\ \forall\, \phi \in C_0^\infty(\Omega_0), \tag{28}$$

and

$$G(\cdot,y) = 0 \text{ on } \partial\Omega_0\,.$$

Furthermore, if $\text{dist}(y, \mathcal{G}\cup\partial\Omega_0) \geqslant \frac{1}{c_1}$ *for some $c_1 > 1$ then*

$$\|G(\cdot,y) - \Gamma(\cdot,y)\|_{H^1(\Omega_0)} \leqslant C, \tag{29}$$

$$\|G(\cdot,y)\|_{H^1(\Omega_0\setminus B_r(y))} \leqslant Cr^{-1/2}, \tag{30}$$

$$\|G(\cdot,y)\|_{L^2(\Omega_0)} \leqslant C, \tag{31}$$

where C depends on $\alpha_0, \beta_0, A, L, \gamma_0, \lambda_1^0$ and on c_1.

The proof of above proposition is similar to the proof of proposition 3.1 in [BFV].

3.4. Unique continuation for the system describing time-harmonic elastic waves

We state a quantitative estimate of unique continuation. We will omit the proof of this estimate since it is a minor modification of the proof of a similar estimate for the Lamé system of elasticity [BFV].

Proposition 3.3. *Let ϵ_1, E_1 and h be positive numbers, $h < h_0$, where h_0 is defined in (24). Let $v \in H^1_{\mathrm{loc}}(\mathcal{K})$ be a solution to*

$$\mathrm{div}(\mathbb{C}\hat{\nabla} v) + \rho\omega^2 v = 0 \text{ in } \mathcal{K},$$

such that

$$\|v\|_{L^\infty(K_0)} \leqslant \epsilon_1$$

and

$$|v(x)| \leqslant E_1(\mathrm{dist}(x, \Sigma_M))^{-1/2} \text{ for every } x \in \mathcal{K}_{h/2}. \tag{32}$$

Then

$$|v(\tilde{x})| \leqslant Cr^{-3/2-\gamma}\epsilon_1^{\tau_r}(E_1 + \epsilon_1)^{1-\tau_r}, \tag{33}$$

where $r \in (0, \frac{1}{C})$, $\tilde{x} = P_M + rn_M$,

$$\tau_r = \tilde{\theta} r^{\delta},$$

and C, δ and $\tilde{\theta}$ with $0 < \tilde{\theta} < 1$ depend on A, L, α_0, β_0, γ_0 and N.

Therefore, if the solution to the system of time-harmonic elastic waves is small in a subdomain of $\mathcal{K}$, and has *a priori* bound (32), then it is also small in $\mathcal{K}$. The above proposition gives a quantitative estimates on how the smallness propagates.

4. Proof of the main result

In this section we prove the main result that consists of showing the uniform continuity for DF and F^{-1}, and establishing a lower bound for DF. These results together with the Fréchet differentiability of F establish theorem 2.7 by proposition 5 of [BV]:

Proposition 4.1. ([*BV*], proposition 5). *Let M_1 and M_2 be positive numbers and $d \in \mathbb{N}$. Let $\mathbf{A}$ and K be an open subset and a compact subset of $\mathbb{R}^d$ respectively. Assume that $K \subset \mathbf{A}$,*

$$\mathrm{dist}\,(K, \mathbb{R}^d \setminus \mathbf{A}) \geqslant M_1, \text{ and } K \subset B_{M_2}(0).$$

Let $\mathcal{B}$ be a Banach space and let $T : \mathbf{A} \to \mathcal{B}$ be such that:

(i) T is Frechét differentiable;
(ii) the Frechét derivative $T' : \mathbf{A} \to \mathcal{L}(\mathbb{R}^d, \mathcal{B})$ is uniformly continuous with a modulus of continuity $\sigma_1(\cdot)$;
(iii) T_K is injective;
(iv) $(T_K)^{-1} : T(K) \to K$ is uniformly continuous with a modulus of continuity $\sigma_2(\cdot)$;
(v) T' is injective in K, namely there is a positive number q_0 such that

$$\min_{x\in K, |h|=1} \|T'(x)[h]\|_{\mathcal{B}} \geqslant q_0;$$

then we have

$$\|x_1 - x_2\|_{\mathbb{R}^d} \leqslant C\|T(x_1) - T(x_2)\|_{\mathcal{B}} \text{ for every } x_1, x_2 \in K,$$

where $C = \max\{\frac{2M_1}{\sigma_2^{-1}(\delta_1)}, \frac{2}{q_0}\}$*, for* $\delta_1 = \frac{1}{2}\min\{\delta_0, M_2\}$ *with* $\delta_0 = \sigma_1^{-1}(\frac{q_0}{2})$*.*

4.1. Injectivity of $F|_{\mathbf{K}}$ *and uniform continuity of* $(F|_{\mathbf{K}})^{-1}$

Let

$$\sigma(t) = \begin{cases} |\log t|^{-\frac{1}{8\delta}} \text{ for } 0 < t < \dfrac{1}{e} \\ t - \dfrac{1}{e} + 1 \text{ for } t \geqslant \dfrac{1}{e} \end{cases} \tag{34}$$

and

$$\sigma_1(t) = (\sigma(t))^{1/5}.$$

Theorem 4.2. *For every* $\underline{L}^1, \underline{L}^2 \in \mathbf{K}$ *the following inequality holds true,*

$$\|\underline{L}^1 - \underline{L}^2\|_\infty \leqslant C_*\sigma_1^N(\|F(\underline{L}^1) - F(\underline{L}^2)\|_\star) \tag{35}$$

where C_* *is a constant depending on* $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0, N$*. Here* $\sigma_1^N(\cdot)$ *is the composition of the function* σ_1 *with itself N times.*

Let $j \in \{1, \ldots, N\}$ be such that

$$d_{D_j}((\mathbb{C}_{\underline{L}^1}, \rho_{\underline{L}^1}), (\mathbb{C}_{\underline{L}^2}, \rho_{\underline{L}^2})) = d_{\Omega_0}((\mathbb{C}_{\underline{L}^1}, \rho_{\underline{L}^1}), (\mathbb{C}_{\underline{L}^2}, \rho_{\underline{L}^2})),$$

and let $D_{j_1}, \ldots, D_{j_M}$ be a chain of domains connecting D_1 to D_j. For the sake of simplicity of notation, set $D_k = D_{j_k}$. Let $\mathcal{W}_k = \text{Int}\,(\cup_{j=0}^k \overline{D}_j)$, $\mathcal{U}_k = \Omega_0 \setminus \mathcal{W}_k$, for $k = 1, \ldots, M-1$. The stiffness tensors $\mathbb{C}_{\underline{L}^1}$ and $\mathbb{C}_{\underline{L}^2}$ are extended as in (22) to all of Ω_0. The densities $\rho_{\underline{L}^1}$ and $\rho_{\underline{L}^2}$ are extended as in (23). We set $\mathbb{C} := \mathbb{C}_{\underline{L}^1}$, $\bar{\mathbb{C}} := \mathbb{C}_{\underline{L}^2}$, $\rho := \rho_{\underline{L}^1}$ and $\bar{\rho} := \rho_{\underline{L}^2}$. Finally, let $\tilde{K}_k = \tilde{K}_h \cap \mathcal{W}_k$ and for $y, z \in \tilde{K}_k$ define the matrix-valued function

$$\mathcal{S}_k(y, z) := \int_{\mathcal{U}_k} ((\mathbb{C} - \bar{\mathbb{C}})\hat{\nabla} G(x, y) : \hat{\nabla}\bar{G}(x, z) - (\rho - \bar{\rho})\omega^2 G(x, y) \cdot \bar{G}(x, z))\mathrm{d}x,$$

the entries of which are given by

$$\mathcal{S}_k^{(p,q)}(y, z) := \int_{\mathcal{U}_k} ((\mathbb{C} - \bar{\mathbb{C}})\hat{\nabla} G^{(p)}(x, y) : \hat{\nabla}\bar{G}^{(q)}(x, z) - (\rho - \bar{\rho})\omega^2 G^{(p)}(x, y) \cdot \bar{G}^{(q)}(x, z))\mathrm{d}x,$$

p, $q = 1, 2, 3$, where $G^{(p)}(\cdot, y)$ and $\bar{G}^{(q)}(, z)$ denote respectively the pth columns and the qth columns of the singular solutions corresponding to $\mathbb{C}, \rho$ and $\bar{\mathbb{C}}, \bar{\rho}$. From (30) we have that

$$|\mathcal{S}_k^{(p,q)}(y, z)| \leqslant C(d(y)d(z))^{-1/2} \text{ for all } y, z \in \tilde{\mathcal{K}}_k,$$

where the constant C depends on the *a priori* parameters only and $d(y) = d(y, \mathcal{U}_k)$ and $d(z) = d(z, \mathcal{U}_k)$.

First, following a similar argument in [BFV], we have the following two propositions:

Proposition 4.3. *For all* $y, z \in \tilde{\mathcal{K}}_k$ *we have that* $\mathcal{S}_k^{(\cdot,q)}(\cdot, z)$*,* $\mathcal{S}_k^{(p,\cdot)}(y, \cdot)$*, belong to* $H^1_{\text{loc}}(\tilde{\mathcal{K}}_k)$ *and for any* $q \in \{1, 2, 3\}$*,*

$$\text{div}(\mathbb{C}\hat{\nabla}\mathcal{S}_k^{(\cdot,q)}(\cdot, z)) + \rho\omega^2\mathcal{S}_k^{(\cdot,q)}(\cdot, z) = 0 \text{ in } \tilde{\mathcal{K}}_k, \tag{36}$$

and for any $p \in \{1, 2, 3\}$,

$$\mathrm{div}(\bar{\mathbb{C}}\hat{\nabla}\mathcal{S}_k^{(p,\cdot)}(y,\cdot)) + \bar{\rho}\omega^2\mathcal{S}_k^{(p,\cdot)}(y,\cdot) = 0 \text{ in } \tilde{\mathcal{K}}_k. \tag{37}$$

Proposition 4.4. *If for a positive* ϵ_0 *and for some* $k \in \{1, \dots, M-1\}$

$$|\mathcal{S}_k(y,z)| \leqslant \epsilon_0 \text{ for every } (y,z) \in K_0 \times K_0, \tag{38}$$

then

$$|\mathcal{S}_k(y_r, z_{\bar{r}})| \leqslant Cr^{-5/2}\bar{r}^{-2}\left(\frac{\epsilon_0}{C_1+\epsilon_0}\right)^{\tau_r\tau_{\bar{r}}}, \tag{39}$$

where $y_r = P_{k+1} + rn_{k+1}$, $z_{\bar{r}} = P_{k+1} + \bar{r}n_{k+1}$, $P_{k+1} \in \Sigma_{k+1}$, $r, \bar{r} \in (0, 1/C)$, $\tau_r = \bar{\theta}r^\delta$, $\tau_{\bar{r}} = \bar{\theta}\bar{r}^\delta$ *and* $C, C_1, \delta, \bar{\theta} \in (0,1)$ *depend on* $A, L, \alpha_0, \beta_0, \gamma_0$ *only.*

We can also prove the following

Proposition 4.5. *If* (38) *holds, then*

$$|\partial_{y_1}\partial_{z_1}\mathcal{S}_k(y_r, z_{\bar{r}})| \leqslant Cr^{-9/2}\bar{r}^{-3}\left(\frac{\epsilon_0}{C_1+\epsilon_0}\right)^{\tau_r\tau_{\bar{r}}}, \tag{40}$$

where $y_r = P_{k+1} + rn_{k+1}$, $z_{\bar{r}} = P_{k+1} + \bar{r}n_{k+1}$, $P_{k+1} \in \Sigma_{k+1}$, $r, \bar{r} \in (0, 1/C)$, $\tau_r = \bar{\theta}r^\delta$, $\tau_{\bar{r}} = \bar{\theta}\bar{r}^\delta$ *and* $C, C_1, \delta, \bar{\theta} \in (0,1)$ *depend on* $A, L, \alpha_0, \beta_0, \gamma_0$ *only.*

We note that, in the above, ∂_{y_1} and ∂_{z_1} denote derivatives in directions lying on the interface Σ_{k+1}.

Proof of proposition 4.5. Fix $z \in K_0$ and consider the function $v(y) := \mathcal{S}^{(\cdot,q)}(y,z)$, for fixed q. By proposition 4.3 we know that v is a solution of

$$\mathrm{div}(\mathbb{C}\hat{\nabla}v(\cdot)) + \rho\omega^2 v(\cdot) = 0 \text{ in } \tilde{\mathcal{K}}_k.$$

Moreover, from proposition 3.2, we get

$$|v(y)| \leqslant C_1 d(y)^{-\frac{1}{2}},\ y \in \tilde{\mathcal{K}}_k,$$

where C_1 depends on $A, L, \alpha_0, \beta_0, \gamma_0, \omega, \lambda_1^0$. Then, applying proposition 3.3 for $\epsilon_1 = \epsilon_0$ and $E_1 = C_1$, we have

$$|v(y_r)| = |\mathcal{S}_k^{(\cdot,q)}(y_r, z)| \leqslant Cr^{-2}\left(\frac{\epsilon_0}{C_1+\epsilon_0}\right)^{\tau_r}$$

for all $y \in B_{r/2}(y_r)$. By the gradient estimate for an elliptic system (see for example [LN]), we obtain

$$|\partial_{y_1}v(y_r)| \leqslant Cr^{-3}\left(\frac{\epsilon_0}{C_1+\epsilon_0}\right)^{\tau_r}.$$

We note that $\partial_{y_1}G(x, y_r) = \partial_{y_1}\Gamma_{k+1}(x, y_r) + \partial_{y_1}w(x, y_r)$, where $\partial_{y_1}w(x, y_r)$ satisfies

$$\begin{cases} \operatorname{div}(\mathbb{C}\hat{\nabla}_x(\partial_{y_1} w(x, y_r))) + \rho\omega^2 \partial_{y_1} w(x, y_r) = \operatorname{div}((\mathbb{C}_b^{k+1} - \mathbb{C})\hat{\nabla}_x(\partial_{y_1}\Gamma_{k+1}(x, y_r))) \\ \qquad\qquad\qquad\qquad\qquad\qquad\qquad -\rho\omega^2 \partial_{y_1}\Gamma_{k+1}(x, y_r) \qquad \text{in } \Omega_0\,, \\ \partial_{y_1} w(x, y_r) = -\partial_{y_1}\Gamma_{k+1}(x, y_r) \qquad\qquad\qquad\qquad\qquad \text{on } \partial\Omega_0\,, \end{cases}$$

where Γ_{k+1} is the biphase fundamental solution for stiffness tensor

$$\mathbb{C}_b^{k+1} = \mathbb{C}_k \chi_{\mathbb{R}^3_+} + \mathbb{C}_{k+1}\chi_{\mathbb{R}^3_-}.$$

Thus $\partial_{y_1} w(\cdot, y_r) \in H^1(\mathcal{U}_k)$ and

$$\|\partial_{y_1} w(\cdot, y_r)\|_{H^1(\mathcal{U}_k)} \leqslant C. \tag{41}$$

Moreover,

$$\begin{aligned} \partial_{y_1} v(y_r) &= \partial_{y_1} S_k^{(\cdot,q)}(y_r, z) \\ &= \int_{\mathcal{U}_k} ((\mathbb{C} - \bar{\mathbb{C}})\hat{\nabla}(\partial_{y_1} G(x, y_r)) : \hat{\nabla}\bar{G}(x, z) \\ &\quad -(\rho - \bar{\rho})\omega^2(\partial_{y_1} G(x, y_r)) \cdot \bar{G}(x, z))\mathrm{d}x, \end{aligned}$$

while

$$\bar{v}(z) = \partial_{y_1} S_k^{(p,\cdot)}(y_r, z),$$

is a solution to

$$\operatorname{div}(\bar{\mathbb{C}}\hat{\nabla} v(\cdot)) + \bar{\rho}\omega^2 v(\cdot) = 0 \text{ in } \tilde{\mathcal{K}}_k,$$

by the same reasoning as in proposition 4.3. By (41) and the estimates,

$$\|\partial_{y_1}\Gamma_{k+1}(\cdot, y)\|_{L^2(\mathbb{R}^3 \setminus B_r(y))} \leqslant Cr^{-1/2}, \tag{42}$$

$$\|\nabla(\partial_{y_1}\Gamma_{k+1}(\cdot, y))\|_{L^2(\mathbb{R}^3 \setminus B_r(y))} \leqslant Cr^{-3/2}, \tag{43}$$

we find that

$$|\bar{v}(z)| \leqslant Cr^{-\frac{3}{2}} d(z)^{-\frac{1}{2}}.$$

Applying proposition 3.3 with $\epsilon_1 = r^{-3}\left(\frac{\epsilon_0}{C_1 + \epsilon_0}\right)^{\tau_r}$ and $E_1 = Cr^{-\frac{3}{2}}$, we have

$$|\bar{v}(z)| \leqslant C\bar{r}^{-2} r^{-\frac{9}{2}} \left(\frac{\epsilon_0}{C_1 + \epsilon_0}\right)^{\tau_r \tau_{\bar{r}}},$$

for all $z \in B_{\bar{r}/2}(z_{\bar{r}})$. Then, again, by the gradient estimate,

$$|\partial_{z_1}\bar{v}(z_{\bar{r}})| \leqslant C\bar{r}^{-3} r^{-\frac{9}{2}} \left(\frac{\epsilon_0}{C_1 + \epsilon_0}\right)^{\tau_r \tau_{\bar{r}}}.$$

Arguing in a similar way, it also follows that

$$
\begin{aligned}
\partial_{z_1}\partial_{y_1}\mathcal{S}_k(y_r, z_{\bar{r}}) &= \partial_{z_1}\bar{v}(z_{\bar{r}}) \\
&= \int_{\mathcal{U}_k} ((\mathbb{C} - \bar{\mathbb{C}})\hat{\nabla}(\partial_{y_1}G(x, y_r)) : \hat{\nabla}(\partial_{z_1}\bar{G}(x, z_{\bar{r}})) \\
&\quad - (\rho - \bar{\rho})\omega^2(\partial_{y_1}G(x, y_r)) \cdot (\partial_{z_1}\bar{G}(x, z_{\bar{r}})))\mathrm{d}x.
\end{aligned}
$$

This completes the proof of (40). □

Proof of theorem 4.2. We follow a walkway and alternate between estimates for Lamé parameters and for the density. Observe that $\|F(\underline{L}^1) - F(\underline{L}^2)\|_\star = \|\Lambda_{\mathbb{C},\rho} - \Lambda_{\bar{\mathbb{C}},\bar{\rho}}\|$. We write

$$
\epsilon := \|F(\underline{L}^1) - F(\underline{L}^2)\|_\star .
$$

Then using (14), we derive that for every $y, z \in K_0$ and for $|l|, |m| = 1$,

$$
\left| \int_\Omega ((\mathbb{C} - \bar{\mathbb{C}})(x)\hat{\nabla}G(x, y)l : \hat{\nabla}\bar{G}(x, z)m - (\rho - \bar{\rho})(x)\omega^2 G(x, y)l \cdot \bar{G}(x, z)m)\mathrm{d}x \right| \leqslant C\epsilon, \tag{44}
$$

where C depends on $\alpha_0, \beta_0, \gamma_0, \omega, A, L$. Let

$$
\delta_k := \max_{0 \leqslant j \leqslant k} \{\max\{|\lambda_j - \bar{\lambda}_j|, |\mu_j - \bar{\mu}_j|, |\rho_j - \bar{\rho}_j|\}\},
$$

where $k \in \{0, 1, \ldots, M\}$.

We will prove that for a suitable, increasing sequence $\{\omega_k(\epsilon)\}_{0 \leqslant k \leqslant M}$ satisfying $\epsilon \leqslant \omega_k(\epsilon)$ for every $k = 0, \ldots, M$ we have

$$
\delta_k \leqslant \omega_k(\epsilon) \Longrightarrow \delta_{k+1} \leqslant \omega_{k+1}(\epsilon), \text{ for every } k = 0, \ldots, M - 1.
$$

Without loss of generality we can choose $\omega_0(\epsilon) = \epsilon$. Suppose now that for some $k = \{1, \ldots, M - 1\}$ we have

$$
\delta_k \leqslant \omega_k(\epsilon). \tag{45}
$$

In the following, we estimate δ_{k+1} by first estimating $|\lambda_{k+1} - \bar{\lambda}_{k+1}|$, $|\mu_{k+1} - \bar{\mu}_{k+1}|$ and then $|\rho_{k+1} - \bar{\rho}_{k+1}|$. Consider

$$
\mathcal{S}_k(y, z) := \int_{\mathcal{U}_k} ((\mathbb{C} - \bar{\mathbb{C}})(x)\hat{\nabla}G(x, y) : \hat{\nabla}\bar{G}(x, z) - (\rho - \bar{\rho})(x)\omega^2 G(x, y) \cdot \bar{G}(x, z))\mathrm{d}x,
$$

and fix $z \in K_0$. From proposition 3.2 and from (44) we get that, for $y, z \in K_0$,

$$
|\mathcal{S}_k(y, z)| \leqslant C(\epsilon + \omega_k(\epsilon)),
$$

where C depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0, \omega$. By (39) and choosing $\bar{r} = cr$ with $c \in [1/4, 1/2]$, we find that there are constants $C_0, \delta \in (0, 1)$ and θ_* depending on A, L, α_0, β_0, γ_0, ω and M, such that for any $r < 1/C_0$ and fixed $l, m \in \mathbb{R}^3$ with $|l| = |m| = 1$,

$$
|\mathcal{S}_k(y_r, z_{\bar{r}})m \cdot l| \leqslant Cr^{-9/2}\varsigma(\omega_k(\epsilon), r), \tag{46}
$$

where

$$
\varsigma(t, s) = \left(\frac{t}{1 + t}\right)^{\theta_* s^{2\delta}}.
$$

We choose $l=m=e_3$ and decompose

$$\mathcal{S}_k(y_r,z_{\bar{r}})e_3\cdot e_3=I_1+I_2, \tag{47}$$

where

$$I_1=\int_{B_{r_1}\cap D_{k+1}}((\mathbb{C}-\bar{\mathbb{C}})(x)\hat{\nabla}G(x,y_r)e_3:\hat{\nabla}\bar{G}(x,z_{\bar{r}})e_3 \\ -(\rho-\bar{\rho})(x)\omega^2G(x,y_r)e_3\cdot\bar{G}(x,z_{\bar{r}})e_3)\mathrm{d}x, \tag{48}$$

$$I_2=\int_{\mathcal{U}_{k+1}\setminus(B_{r_1}\cap D_{k+1})}((\mathbb{C}-\bar{\mathbb{C}})(x)\hat{\nabla}G(x,y_r)e_3:\hat{\nabla}\bar{G}(x,z_{\bar{r}})e_3 \\ -(\rho-\bar{\rho})(x)\omega^2G(x,y_r)e_3\cdot\bar{G}(x,z_{\bar{r}})e_3)\mathrm{d}x, \tag{49}$$

with $r_1=\frac{1}{4LC_L}$. Then, from proposition 3.2, we derive immediately that

$$|I_2|\leqslant C. \tag{50}$$

By (31), we have

$$\left|\int_{B_{r_1}\cap D_{k+1}}(\rho-\bar{\rho})(x)\omega^2G(x,y_r)e_3\cdot\bar{G}(x,z_{\bar{r}})e_3\mathrm{d}x\right|\leqslant C,$$

where C depends on $A,L,\alpha_0,\beta_0,\gamma_0,\lambda_1^0$. Using (29) and (30), we get

$$|I_1|\geqslant\left|\int_{B_{r_1}\cap D_{k+1}}(\mathbb{C}_b^{k+1}-\bar{\mathbb{C}}_b^{k+1})(x)\hat{\nabla}\Gamma_{k+1}(x,y_r)e_3:\hat{\nabla}\bar{\Gamma}_{k+1}(x,z_{\bar{r}})e_3\mathrm{d}x\right| \\ -C\left(\frac{1}{\sqrt{r}}+1\right), \tag{51}$$

where Γ_{k+1} and $\bar{\Gamma}_{k+1}$ are the biphase fundamental solutions introduced in section 3.3 corresponding to the stiffness tensors $\mathbb{C}_b^{k+1}$ and $\bar{\mathbb{C}}_b^{k+1}$ given by

$$\mathbb{C}_b^{k+1}=\mathbb{C}_k\chi_{\mathbb{R}^3_+}+\mathbb{C}_{k+1}\chi_{\mathbb{R}^3_-},$$

$$\bar{\mathbb{C}}_b^{k+1}=\bar{\mathbb{C}}_k\chi_{\mathbb{R}^3_+}+\bar{\mathbb{C}}_{k+1}\chi_{\mathbb{R}^3_-},$$

up to a rigid coordinate transformation that maps the flat part of Σ_{k+1} into $x_3=0$. Furthermore by (46), (47) and (50) we obtain

$$|I_1|\leqslant C(r^{-9/2}\varsigma(\omega_k(\epsilon),r)+1), \tag{52}$$

where C depends on $A,L,\alpha_0,\beta_0,\gamma_0,\lambda_1^0$. Hence, by (51) and (52) and by performing the change of variables $x=rx'$ in the integral, we get

$$\left|\int_{B^-_{r_1/r}}(\mathbb{C}_b^{k+1}-\bar{\mathbb{C}}_b^{k+1})(x')\hat{\nabla}\Gamma_{k+1}(x',e_3)e_3:\hat{\nabla}\bar{\Gamma}_{k+1}(x',ce_3)e_3\mathrm{d}x'\right|\leqslant\delta_0(r), \tag{53}$$

where

$$\delta_0(r) = C\,[r^{-7/2}\varsigma(\omega_k(\epsilon), r) + r^{1/2}]\,.$$

We then follow the procedure of [BFV] pp 632–4, and obtain

$$|\lambda_{k+1} - \bar{\lambda}_{k+1}| \leqslant C\sigma(\omega_k(\epsilon)), \quad |\mu_{k+1} - \bar{\mu}_{k+1}| \leqslant C\sigma(\omega_k(\epsilon)). \tag{54}$$

Next, we estimate $|\rho_{k+1} - \bar{\rho}_{k+1}|$. By proposition 4.5, there are constants $C_0, \delta \in (0, 1)$ and θ_* depending on $A, L, \alpha_0, \beta_0, \gamma_0, \omega$ and, increasingly, on M, such that for any $r < 1/C_0$ and fixed $l, m \in \mathbb{R}^3$ such that $|l| = |m| = 1$,

$$|\partial_{y_1}\partial_{z_1}\mathcal{S}_k(y_r, y_r)m \cdot l| \leqslant Cr^{-15/2}\varsigma(\omega_k(\epsilon), r). \tag{55}$$

We choose $l = m = e_3$, again, and decompose

$$\partial_{y_1}\partial_{z_1}\mathcal{S}_k(y_r, y_r)e_3 \cdot e_3 = J_1 + J_2, \tag{56}$$

where

$$\begin{aligned} J_1 = \int_{B_{r_1}\cap D_{k+1}} &((\mathbb{C} - \bar{\mathbb{C}})(x)\hat{\nabla}\ (\partial_{y_1}G(x, y_r))e_3 : \hat{\nabla}(\partial_{z_1}\bar{G}(x, y_r))e_3 \\ &- (\rho - \bar{\rho})(x)\omega^2(\partial_{y_1}G(x, y_r))e_3 \cdot (\partial_{z_1}\bar{G}(x, y_r))e_3)\mathrm{d}x, \end{aligned} \tag{57}$$

$$\begin{aligned} J_2 = \int_{\mathcal{U}_{k+1}\setminus(B_{r_1}\cap D_{k+1})} &((\mathbb{C} - \bar{\mathbb{C}})(x)\hat{\nabla}\ (\partial_{y_1}G(x, y_r))e_3 : \hat{\nabla}(\partial_{z_1}\bar{G}(x, y_r))e_3 \\ &- (\rho - \bar{\rho})(x)\omega^2(\partial_{y_1}G(x, y_r))e_3 \cdot (\partial_{z_1}\bar{G}(x, y_r))e_3)\mathrm{d}x. \end{aligned} \tag{58}$$

Then, with (41)–(43) we derive that

$$|J_2| \leqslant C. \tag{59}$$

By estimates (41)–(43), and using that $|\lambda_k - \bar{\lambda}_k| \leqslant C\omega_k(\epsilon)$, $|\mu_k - \bar{\mu}_k| \leqslant C\omega_k(\epsilon)$, $|\lambda_{k+1} - \bar{\lambda}_{k+1}| \leqslant C\sigma(\omega_k(\epsilon))$ and $|\mu_{k+1} - \bar{\mu}_{k+1}| \leqslant C\sigma(\omega_k(\epsilon))$, we get

$$\begin{aligned} |J_1| \geqslant &\left| \int_{B_{r_1}\cap D_{k+1}} (\rho_{k+1} - \bar{\rho}_{k+1})\frac{\partial}{\partial y_1}\Gamma_{k+1}(x, y_r)e_3 \cdot \frac{\partial}{\partial y_1}\Gamma_{k+1}(x, y_r)e_3\mathrm{d}x \right| \\ &- C\left(\frac{1}{\sqrt{r}} + \frac{\sigma(\omega_k(\epsilon))}{r^3}\right) \\ \geqslant &\ |\rho_{k+1} - \bar{\rho}_{k+1}| \int_{B_{r_1}\cap D_{k+1}} \left|\frac{\partial}{\partial y_1}\Gamma_{k+1}(x, y_r)e_3\right|^2 \mathrm{d}x - C\left(\frac{1}{\sqrt{r}} + \frac{\sigma(\omega_k(\epsilon))}{r^3}\right), \end{aligned} \tag{60}$$

where we have used that

$$\int_{B_{r_1}\cap D_{k+1}} \left|\frac{\partial}{\partial y_1}\Gamma_{k+1}(x, y_r)e_3\right| \left|\frac{\partial}{\partial y_1}\Gamma_{k+1}(x, y_r)e_3 - \frac{\partial}{\partial y_1}\bar{\Gamma}_{k+1}(x, y_r)e_3\right| \mathrm{d}x \leqslant C\frac{\sigma(\omega_k(\epsilon))}{r}.$$

Furthermore, by (55),(56) and (59) we obtain

$$|J_1| \leqslant C(r^{-15/2}\varsigma(\omega_k(\epsilon), r) + 1). \tag{61}$$

By (60) and by performing the change of variables $x = rx'$ in the integral, we have

$$r^{-1}|\rho_{k+1} - \bar{\rho}_{k+1}| \int_{B^-_{r_1/r}} \left| \frac{\partial}{\partial y_1}\Gamma_{k+1}(x', e_3)e_3 \right|^2 \mathrm{d}x' \leqslant C\left((r^{-15/2}\varsigma(\omega_k(\epsilon), r) + 1) + \frac{1}{\sqrt{r}} + \frac{\sigma(\omega_k(\epsilon))}{r^3}\right).$$

Since $r_1/r \geqslant C/4LC_L$ when $r \in (0, 1/C)$, we have

$$\int_{B^-_{r_1/r}} \left| \frac{\partial}{\partial y_1}\Gamma_{k+1}(x', e_3)e_3 \right|^2 \mathrm{d}x' \geqslant \int_{B^-_{C/4LC_L}} \left| \frac{\partial}{\partial y_1}\Gamma_{k+1}(x', e_3)e_3 \right|^2 \mathrm{d}x' \geqslant C,$$

for some positive C. Then

$$|\rho_{k+1} - \bar{\rho}_{k+1}|r^{-1} \leqslant C\left((r^{-15/2}\varsigma(\omega_k(\epsilon), r) + 1) + \frac{1}{\sqrt{r}} + \frac{\sigma(\omega_k(\epsilon))}{r^3}\right),$$

and thus

$$|\rho_{k+1} - \bar{\rho}_{k+1}| \leqslant \delta_1(r), \tag{62}$$

where

$$\delta_1(r) = C\left[r^{-13/2}\varsigma(\omega_k(\epsilon), r) + \sqrt{r} + \frac{\sigma(\omega_k(\epsilon))}{r^2}\right].$$

If $\omega_k(\epsilon) < 1/e$, we choose

$$r = \frac{|\sigma(\omega_k(\epsilon))|^{2/5}}{C},$$

and then

$$|\rho_{k+1} - \bar{\rho}_{k+1}| \leqslant C|\sigma(\omega_k(\epsilon))|^{1/5}. \tag{63}$$

Otherwise, if $\omega_k(\epsilon) \geqslant 1/e$, since $|\rho_{k+1} - \bar{\rho}_{k+1}|$ is bounded, we get (63) trivially. By (54) and (63), we follow the weakest estimate to get

$$\delta_{k+1} \leqslant \omega_{k+1}(\epsilon) := C\sigma_1(\omega_k(\epsilon)).$$

Following the way of alternatingly estimating $|\lambda - \bar{\lambda}|, |\mu - \bar{\mu}|$ and $|\rho - \bar{\rho}|$ along the walkay $D_1, D_2, \ldots, D_M$, and recalling that $\omega_0(\epsilon) = \epsilon$, we get (35). □

The uniqueness statement in theorem 2.7 is an immediate consequence of the proposition above.

4.2. Injectivity of $DF(\underline{L})$ and estimate from below of $DF|_{\mathbf{K}}$

Proposition 4.6. *Let*

$$q_0 := \min\{\|DF(\underline{L})[\underline{H}]\|_\star \mid \underline{L} \in \mathbf{K}, \underline{H} \in \mathbb{R}^{3N}, \|\underline{H}\|_\infty = 1\};$$

we have

$$(\sigma_1^N)^{-1}(1/C_\star) \leqslant q_0, \tag{64}$$

where $C_\star > 1$ depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0$ and N only.

Proof. By the definition of q_0 there exists an $\underline{L}_0 \in \mathbf{K}$ and

$$\underline{H}_0 = (h_{0,1}, \dots, h_{0,N}, k_{0,1}, \dots, k_{0,N}, l_{0,1}, \dots, l_{0,N}), \; \|\underline{H}_0\|_\infty = 1,$$

such that

$$\|DF(\underline{L}_0)[\underline{H}_0]\|_\star = q_0. \tag{65}$$

Therefore, by (15), (65), we have

$$\left| \int_\Omega \mathbb{H}(x)(\hat{\nabla} G(x,y)l : \hat{\nabla} G(x,z)m - h(x)\omega^2 G(x,y)l \cdot G(x,z)m)\mathrm{d}x \right| \leqslant Cq_0 \tag{66}$$

for every $y, z \in \mathcal{K}_0$, where C depends on $\alpha_0, \beta_0, \gamma_0, \omega, A, L$, $\mathbb{H} = \mathbb{C}_{\underline{H}_0}$, $h = \rho_{\underline{H}_0}$ and $G(\cdot, y)$ denotes the singular solution corresponding to $\mathbb{C}_{\underline{L}}, \rho_{\underline{L}}$. From now on the vector

$$(0, h_{0,1}, \dots, h_{0,N}, 0, k_{0,1}, \dots, k_{0,N}, 0, l_{0,1}, \dots, l_{0,N}),$$

will still be denoted by $\underline{H}_0$.

We fix $j \in \{1, \dots, N\}$ and let $D_{j_1}, \dots, D_{j_M}$ be a chain of domains connecting D_1 to D_j, where

$$\max\{|h_{0,j}|, |k_{0,j}|, |l_{0,j}|\} = \|\underline{H}_0\|_\infty = 1.$$

Now, let

$$\eta_i := \max_{0 \leqslant j \leqslant i}\{\max\{|h_{0,j}|, |k_{0,j}|, |l_{0,j}|\}\},$$

where $i \in \{0, 1, \dots, M\}$.

We will prove that for a suitable increasing sequence $\{\omega_i(q_0)\}_{0 \leqslant i \leqslant M}$ satisfying $\epsilon \leqslant \omega_i(q_0)$ for every $k = 0, \dots, M$, we have

$$\delta_k \leqslant \omega_i(q_0) \Longrightarrow \delta_{i+1} \leqslant \omega_{k+1}(q_0) \text{ for every } i = 0, \dots, M-1.$$

Without loss of generality we can choose $\omega_0(q_0) = q_0$. Suppose now that for some $i = \{1, \dots, M-1\}$ we obtain (65). Let $\mathcal{Y}_i(y,z) = \{\mathcal{Y}_i^{(p,q)}(y,z)\}_{1 \leqslant p,q \leqslant 3}$ be the matrix valued function the elements of which are given by

$$\mathcal{Y}_i^{(p,q)}(y,z) := \int_{\mathcal{U}_i} (\mathbb{H}(x)\hat{\nabla} G^{(p)}(x,y) : \hat{\nabla} G^{(q)}(x,z) - h(x)\omega^2 G^{(p)}(x,y) \cdot G^{(q)}(x,z))\mathrm{d}x,$$

with $z \in K_0$ fixed. From proposition 3.2 and from (44) we get that, for $y, z \in K_0$,

$$|\mathcal{Y}_i(y, z)| \leqslant C(q_0 + \omega_i(q_0)),$$

where C depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0$. Choosing $\bar{r} = cr$ with $c \in [1/4, 1/2]$, as in proposition 4.4, we have that there exists a constant C_2 such that for every $r \in (0, 1/C_2)$,

$$|\mathcal{Y}_i(y_r, z_{\bar{r}})| \leqslant Cr^{-9/2}\varsigma(\omega_i(q_0, r)), \tag{67}$$

where

$$\varsigma(t, s) = \left(\frac{t}{1+t}\right)^{\theta_* s^{2\delta}}.$$

We choose $l = m = e_3$, again, and decompose

$$\mathcal{Y}_k(y_r, z_{\bar{r}})e_3 \cdot e_3 = I_1 + I_2, \tag{68}$$

where

$$I_1 = \int_{B_{r_1} \cap D_{i+1}} (\mathbb{H}(x)\hat{\nabla} G(x, y_r)e_3 : \hat{\nabla} G(x, z_{\bar{r}})e_3 - h(x)\omega^2 \bar{G}(x, y_r)e_3 \cdot G(x, z_{\bar{r}})e_3)\mathrm{d}x, \tag{69}$$

$$\begin{aligned} I_2 = \int_{\mathcal{U}_{i+1} \backslash (B_{r_1} \cap D_{i+1})} & (\mathbb{H}(x)\hat{\nabla} G(x, y_r)e_3 : \hat{\nabla} G(x, z_{\bar{r}})e_3 \\ & - h(x)\omega^2 G(x, y_r)e_3 \cdot G(x, z_{\bar{r}})e_3)\mathrm{d}x, \end{aligned} \tag{70}$$

and $r_1 = \frac{1}{4LC_L}$. Then, from proposition 3.2, we derive that

$$|I_2| \leqslant C. \tag{71}$$

Using (31), we find that

$$\left| \int_{B_{r_1} \cap D_{k+1}} h(x)\omega^2 G(x, y_r)e_3 \cdot G(x, z_{\bar{r}})e_3 \mathrm{d}x \right| \leqslant C,$$

where C depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0$. Then, by (29) and (30) we get

$$|I_1| \geqslant \left| \int_{B_{r_1} \cap D_{i+1}} \mathbb{H}(x)\hat{\nabla}\Gamma_{i+1}(x, y_r)e_3 : \hat{\nabla}\Gamma_{i+1}(x, z_{\bar{r}})e_3 \mathrm{d}x \right| - C\left(\frac{1}{\sqrt{r}} + 1\right). \tag{72}$$

With (67), (68) and (71) we obtain

$$|I_1| \leqslant C(r^{-9/2}\varsigma(\omega_i(q_0), r) + 1), \tag{73}$$

where C depends on $A, L, \alpha_0, \beta_0, \gamma_0, \lambda_1^0$. Following the procedure of [BFV] pp. 635-637, we get

$$|h_{0,i+1}| \leqslant C\sigma(\omega_i(q_0)), \quad |k_{0,i+1}| \leqslant C\sigma(\omega_i(q_0)). \tag{74}$$

Similar to proposition 4.5, we find that there are constants $C_2, \delta \in (0, 1)$ and θ_* depending on $A, L, \alpha_0, \beta_0, \gamma_0, \omega$ and, increasingly, on M, such that for any $r < 1/C_2$

$$|\partial_{y_1}\partial_{z_1}\mathcal{Y}_i(y_r, y_r)e_3 \cdot e_3| \leqslant Cr^{-15/2}\varsigma(\omega_i(q_0, r)). \tag{75}$$

We decompose

$$\partial_{y_1}\partial_{z_1}\mathcal{Y}_i(y_r, y_r)e_3 \cdot e_3 = J_1 + J_2, \tag{76}$$

where

$$J_1 = \int_{B_{r_1}\cap D_{i+1}} (\mathbb{H}(x)\hat{\nabla}(\partial_{y_1}G(x, y_r))e_3 : \hat{\nabla}\ (\partial_{z_1}G(x, y_r))e_3 \\ - h(x)\omega^2(\partial_{y_1}G(x, y_r))e_3 \cdot (\partial_{z_1}G(x, y_r))e_3)\mathrm{d}x, \tag{77}$$

$$J_2 = \int_{\mathcal{U}_{i+1}\setminus(B_{r_1}\cap D_{i+1})} (\mathbb{H}(x)\hat{\nabla}(\partial_{y_1}G(x, y_r))e_3 : \hat{\nabla}\ (\partial_{z_1}G(x, y_r))e_3 \\ - h(x)\omega^2(\partial_{y_1}G(x, y_r))e_3 \cdot (\partial_{z_1}G(x, y_r))e_3)\mathrm{d}x. \tag{78}$$

Using (41)–(43) and (74), we get

$$|J_2| \leqslant C \tag{79}$$

and

$$|J_1| \geqslant \left| \int_{B_{r_1}\cap D_{i+1}} l_{0,i+1}\frac{\partial}{\partial y_1}\Gamma_{i+1}(x, y_r)e_3 \cdot \frac{\partial}{\partial y_1}\Gamma_{i+1}(x, y_r)e_3\mathrm{d}x \right| - C\left(\frac{1}{\sqrt{r}} + \frac{\sigma(\omega_i(\epsilon))}{r^3}\right) \\ = \quad |l_{0,i+1}| \int_{B_{r_1}\cap D_{i+1}} \left|\frac{\partial}{\partial y_1}\Gamma_{i+1}(x, y_r)e_3\right|^2 \mathrm{d}x - C\left(\frac{1}{\sqrt{r}} + \frac{\sigma(\omega_i(q_0))}{r^3}\right). \tag{80}$$

Furthermore by (75), (76) and (79), we obtain

$$|J_1| \leqslant C(r^{-15/2}\varsigma((\omega_i(q_0)), r) + 1). \tag{81}$$

Hence, by (80) and upon performing the change of variables $x = rx'$ in the integral, we obtain

$$r^{-1}|l_{0,i+1}| \int_{B^-_{r_1/r}} \left|\frac{\partial}{\partial y_1}\Gamma_{i+1}(x', e_3)e_3\right|^2 \mathrm{d}x' \\ \leqslant C\left((r^{-15/2}\varsigma(\omega_i(q_0), r) + 1) + \frac{1}{\sqrt{r}} + \frac{\sigma(\omega_i(q_0))}{r^3}\right).$$

Since $r_1/r \geqslant C/4LC_L$ when $r \in (0, 1/C)$, we have

$$\int_{B^-_{r_1/r}} \left|\frac{\partial}{\partial y_1}\Gamma_{i+1}(x', e_3)e_3\right|^2 \mathrm{d}x' \geqslant \int_{B^-_{C/4LC_L}} \left|\frac{\partial}{\partial y_1}\Gamma_{i+1}(x', e_3)e_3\right|^2 \mathrm{d}x' \geqslant C.$$

Then

$$|l_{0,i+1}|r^{-1} \leqslant C\left((r^{-15/2}\varsigma((\omega_i(q_0)), r) + 1) + \frac{1}{\sqrt{r}} + \frac{\sigma(\omega_i(q_0))}{r^3}\right),$$

and thus

$$|l_{0,i+1}| \leqslant \delta_1(r), \tag{82}$$

where

$$\delta_1(r) = C\left[r^{-13/2}\varsigma(\omega_i(q_0), r) + \sqrt{r} + \frac{\sigma(\omega_i(q_0))}{r^2}\right].$$

If $\omega_i(q_0) < 1/e$, we choose

$$r = \frac{|\sigma(\omega_i(q_0))|^{2/5}}{C}$$

so that

$$|l_{0,i+1}| \leqslant C|\sigma(\omega_i(q_0))|^{1/5}. \tag{83}$$

Otherwise, if $\omega_i(q_0) \geqslant 1/e$, because $|l_{0,i+1}|$ is bounded, we get (83) trivially. Then, by (74) and (83) we get

$$\eta_{i+1} \leqslant \omega_{i+1}(q_0) := C\sigma_1(\omega_i(q_0)).$$

Finally, by alternating the estimates for $|\lambda - \bar{\lambda}|$, $|\mu - \bar{\mu}|$ and $|\rho - \bar{\rho}|$, we get

$$1 = \eta_M \leqslant C\sigma_1^M(q_0) \leqslant C\sigma_1^N(q_0),$$

and the statement follows. □

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